

**Definition 15.** (*Finite Element space*)  
 Given discretization  $\Omega_h$  of  $\Omega \subset \mathbb{R}^d$ , the space

$$V_h^p(\mathcal{T}) = \{u \in H^1 : \text{für alle } T \in \mathcal{T} : u|_T \in \mathbb{P}_p\}$$

is called the *conforming Finite Element space of order  $p$* .

From the example of Courant a general method for  $\mathbb{R}^d$ ,  $d = 1, 2, 3$ , can be derived.

### 5.11.3 Finite Elements in $\mathbb{R}^1$

Let us consider the Helmholtz equation

$$-\Delta u + u = f \text{ with } a(u, v) = \int_{\Omega} (\nabla u \nabla v + uv) dx \quad (5.39)$$

Let  $\mathcal{N} = \{a = x_0, x_1, x_2, \dots, x_{N+1} = b\}$  be the discretization of domain  $[a, b]$  with element width  $h_i = x_{i+1} - x_i$  and local basis functions  $\varphi_i$  with the following property:

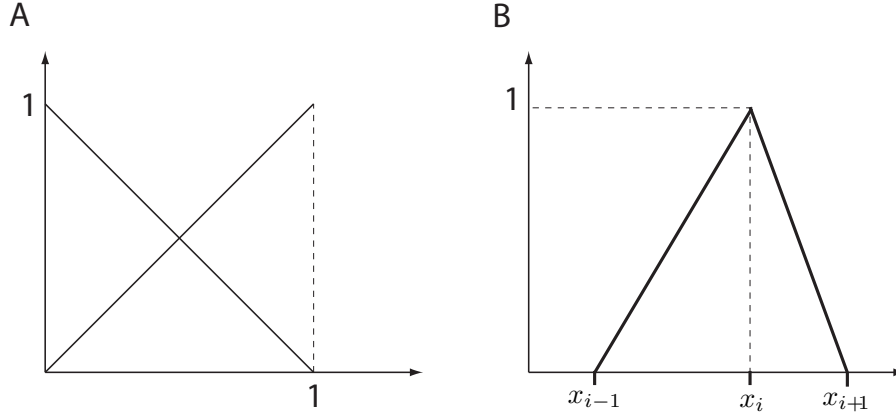
$$\varphi_i(x_j) = \delta_{ij}$$

Then the solution  $u_h$  has the form

$$u_h = \sum_{i=1}^N a_i \varphi_i.$$

#### Computing $u_h$ on a reference interval

The basis functions  $\varphi_i$  can be expressed by form functions  $\Phi_i$ . Using a bijective, affine, transformational mapping, the interval  $[x_i, x_{i+1}]$  can be mapped to the reference interval  $[0, 1]$ . The solution  $u_i$  can then be computed on the interval  $[0, 1]$ , expressed by the form functions and later transferred back to the original element.



### Transformational mapping

Let  $I_i = [x_i, x_{i+1}]$  and  $\xi \in [0, 1]$ . Then the following functions can be defined:

$$\begin{aligned}
 x_{I_i} : [0, 1] &\longrightarrow I_i \\
 \xi &\longmapsto x_i + h_i \xi \\
 \xi_{I_i} : I_i &\longrightarrow [0, 1] \\
 x &\longmapsto \frac{(x - x_i)}{h_i}
 \end{aligned}$$

This defines a bijection between  $[0, 1]$  and  $[x_i, x_{i+1}]$ . On the reference interval the solution can be expressed as:

$$u_h(\xi) = \alpha_1 + \alpha_2 \xi.$$

Here,  $u_i = u_h(0) = \alpha_1$  and  $u_{i+1} = u_h(1) = \alpha_1 + \alpha_2$ .

$$\begin{aligned}
 \Rightarrow u_h(\xi) &= \alpha_1 + \alpha_2 \xi = u_i + (u_{i+1} - u_i) \xi \\
 &= (1 - \xi) u_i + \xi u_{i+1} = u_i \Phi_1(\xi) + u_{i+1} \Phi_2(\xi)
 \end{aligned}$$

**Note 8.** The following property applies to the form functions:

$$\forall \xi \in [0, 1] : \Phi_1(\xi) + \Phi_2(\xi) = 1$$

For the basis functions we get:

$$\varphi_i(x) = \begin{cases} \Phi_2(\xi(x)) & x \in I_{i-1} \\ \Phi_1(\xi(x)) & x \in I_i \\ 0 & \text{else} \end{cases}$$

Now we can compute the matrix entries of the linear system in terms of the form functions.

### Computing the system matrix entries

We need to compute  $a(\varphi_i, \varphi_j)$  of system matrix  $A$  (in our case for the Helmholtz problem):

$$a(\varphi_i, \varphi_j) = \sum_{k=1}^N \int_{I_k} \nabla \varphi_i(x) \nabla \varphi_j(x) + \varphi_i(x) \varphi_j(x) dx.$$

The functions  $\varphi_i, \varphi_j$  can be replaced by  $\Phi_n$  and  $\Phi_m$  ( $n, m \in \{1, 2\}$ ).

$$\begin{aligned} \Rightarrow (A_{I_k})_{nm} &= \int_{I_k} \nabla_x \Phi_n(\xi(x)) \nabla_x \Phi_m(\xi(x)) + \Phi_n(\xi(x)) \Phi_m(\xi(x)) dx \\ &= (x_{k+1} - x_k) \int_0^1 \nabla_\xi \Phi_n(\xi) \xi'(x(\xi)) \xi'(x(\xi)) \nabla_\xi \Phi_m(\xi) + \Phi_n \Phi_m d\xi \\ &= h_k \int_0^1 \frac{1}{h_k^2} \nabla \Phi_n \nabla \Phi_m + \Phi_n \Phi_m d\xi \end{aligned}$$

We see that the matrix is comprised of  $2 \times 2$  sub-matrices  $A_{I_k}$ .

The matrix entries of  $A_{I_k}$  are:

$$\begin{aligned} (A_{I_k})_{11} &= \int_0^1 \frac{1}{h_k} \nabla \Phi_1 \nabla \Phi_1 + \Phi_1 \Phi_1 \cdot h_k d\xi \\ &= \int_0^1 \frac{1}{h_k} + h_k (1 - \xi)^2 d\xi = \frac{1}{h_k} + \frac{1}{3} h_k \\ (A_{I_k})_{12} &= \int_0^1 \frac{1}{h_k} \nabla \Phi_1 \nabla \Phi_2 + \Phi_1 \Phi_2 \cdot h_k d\xi \\ &= \int_0^1 -\frac{1}{h_k} + \xi(1 - \xi) \cdot h_k d\xi \\ &= -\frac{1}{h_k} + \frac{1}{6} h_k \\ (A_{I_k})_{21} &= -\frac{1}{h_k} + \frac{1}{6} h_k \\ (A_{I_k})_{22} &= \frac{1}{h_k} + \frac{1}{3} h_k \end{aligned}$$

$$\Rightarrow A_{I_k} = \frac{1}{h_k} \begin{bmatrix} 1 & -1 \\ -1 & 1 \end{bmatrix} + h_k \begin{bmatrix} 1/3 & 1/6 \\ 1/6 & 1/3 \end{bmatrix}$$

### Quadratic elements

Instead of linear approximation, we can use quadratic approximation (still low order) which should increase the approximation quality. This leads to

$$u_h(\xi) = \alpha_1 + \alpha_2 \xi + \alpha_3 \xi^2 \text{ on } I = [0, 1].$$